THE SEARCH FOR A PERFECT FLUID AND TRANSPORT WITHOUT QUASIPARTICLES

work done together with



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Scale invariance and non-relativistic CFT's Nishida/Son '07

dilatation operator generates $x \to e^{\lambda} x$ and $t \to e^{2\lambda} t$

 \hat{H} is scale-invariant if $i\left[\hat{H},\hat{D}\right] = 2\,\hat{H} = \frac{d}{dt}\int x_i\hat{j}_i$

momentum balance $\partial_t \hat{j}_i = -\partial_j \hat{\Pi}_{ij} \rightarrow$

$$\frac{d}{dt}\hat{D} = 2\int \hat{\epsilon} = \int x_i \partial_t \hat{j}_i = -\int x_i \partial_j \hat{\Pi}_{ij} = \int \hat{\Pi}_{ii}$$

scale invariance implies $|2\epsilon = \prod_{ii}|$ trace of stress tensor

conformal transformation $x \to x/(1 + \lambda t)$ and $t \to t/(1 + \lambda t)$

$$\hat{H}$$
 is conformal-invariant if $i\left[\hat{H},\hat{C}\right] = \hat{D} = \frac{d}{dt}\int x^2 \hat{n}/2 \rightarrow \hat{D}$

breathing mode in a trap at $\omega_B = 2\omega_{\text{trap}}$ Castin/Werner '06

The unitary gas as a quantum critical point Nikolic/Sachdev '07

action of free Fermions
$$S_0 = \int d\tau \int d^d x \left[\psi^* \partial_\tau \psi + \frac{\hbar^2}{2m} |\nabla \psi|^2 \right]$$

is invariant under $x \to x e^{-l}, \tau \to \tau e^{-zl}$ and $\psi \to \psi e^{dl/2} \to$

dyn. exponent z = 2 and dim $[\psi] = d/2$

chemical potential $\mathcal{L}_{\mu} = -\mu |\psi|^2$ scales like $\mu \to \mu e^{2l} \to \dim [\mu] = 2$

add zero range interaction $\mathcal{L}_{int} = u_0 \psi^*_{\uparrow} \psi^*_{\downarrow} \psi_{\downarrow} \psi_{\uparrow} \rightarrow \dim [u_0] = 2 - d$



scaling of the dim.less coupling

$$u = 2mS_d \Lambda^{d-2} u_0 \text{ under } \Lambda \to \Lambda e^{-l}$$

$$du/dl = \epsilon u - u^2/2$$
 with $\epsilon = 2 - d$

unitary Fermions \rightarrow zero density QCP



single length scale $\lambda_T \rightarrow \text{density } n\lambda_T^3 \simeq 3.1 \text{ vanishes} \sim T^{3/2}$ universal amplitude ratio e.g. $s(T) \sim k_B \cdot \lambda_T^{-3}$ and $\eta(T) \sim \hbar \cdot \lambda_T^{-3}$

 $\rightarrow \eta/s \simeq 0.7\hbar/k_B$ characterizes a strongly coupled QFT

dimensional analysis suffices if there is no anomalous dimension !

The unitary gas as a 'perfect fluid'





Definition A fluid is perfect if

$$\frac{\eta}{s} \equiv \frac{\hbar}{4\pi k_B}$$

Starinets '05 (SSYM) All known fluids have

$$\frac{\eta}{s} \ge \frac{\hbar}{4\pi k_B} \,\, !!$$



momentum balance $\partial_t(\rho v_i) + \partial_i \Pi_{ij} = 0$

$$\Pi_{ij} = p\delta_{ij} + \rho v_i v_j - \eta \left(\partial_i v_j + \partial_j v_i - \frac{2}{3} \delta_{ij} \cdot \partial_k v_k \right) - \zeta \delta_{ij} \cdot \partial_k v_k$$

positivity $\eta \ge 0$ and $\zeta \ge 0$ due to $dS/dt \ge 0$

liquids thermally activated $\rightarrow \eta(T)$ grows as $T \downarrow$

gases
$$\eta = \frac{1}{3}mn \langle v \rangle \ell \simeq \sqrt{mk_BT} / \sigma(T)$$
 grows as $T \uparrow$

A quantum limit on the viscosity ?

mean free path $\ell \gtrsim n^{-1/3}$ average velocity $\langle v \rangle \gtrsim \frac{\hbar}{m} n^{1/3}$

gives $\eta \ge \alpha_{\eta} \cdot \hbar n$ e.g. $\alpha_{\eta} \simeq 0.5$ for ⁴He at 2K



Cao ... Science **331** (2011) and Sommer ... Nature **472** (2011)

shear viscosity of the unitary gas

Boltzmann-limit $\eta(T \gg T_F) = 2.8 \hbar n (T/T_F)^{3/2} = 4.2 \frac{\hbar}{\lambda_T^3}$ (density drops out!), well defined quasipart. $\hbar/\tau_\eta \ll k_B T$ **superfluid** below $T_c \simeq 0.16 T_F$ has **finite** viscosity due to a) phonon interactions: $\eta(T) \sim T^{-5}$ as $T \ll T_c$ Rupak/Schäfer '07 b) fermionic qp's: $\eta(T) \rightarrow \text{const} \text{ as } T \rightarrow 0$ Pethick/Smith '75

minimum is observed in $^{4}{\rm He}$ just below T_{λ}



Viscosity in linear response: Kubo formula

• viscosity from stress correlations (cf. hydrodynamics):

$$\eta(\omega) = \frac{1}{\omega} \operatorname{Re} \int_0^\infty dt \, e^{i\omega t} \int d^3 x \left\langle \left[\hat{\Pi}_{xy}(\boldsymbol{x}, t), \hat{\Pi}_{xy}(0, 0) \right] \right\rangle$$

with stress tensor $\hat{\Pi}_{xy} = \sum_{\mathbf{p}, \sigma} \frac{p_x p_y}{m} c^{\dagger}_{\mathbf{p}\sigma} c_{\mathbf{p}\sigma} \quad \text{(cf. Newton } \frac{\partial v_x}{\partial y}\text{)}$

• correlation function (Kubo formula): Enss, Haussmann & Zwerger Ann. Phys. 2011



- transport via fermions and bosonic molecules: very efficient description, satisfies conservation laws, scale invariance and Tan relations Enss PRA 2012
- assumes no quasiparticles: beyond Boltzmann kinetic theory, works near Tc

Dynamic shear viscosity



Enss, Haussmann & Zwerger 2011; Enss 2013; cf. Taylor & Randeria 2010



Shear viscosity/entropy of the unitary Fermi gas

Enss, Haussmann & Zwerger 2011

Spin transport with ultracold gases

• experiment: spin-polarized clouds in harmonic trap



• strongly interacting gas [movie courtesy Martin Zwierlein]:



bounce!

A.T. Sommer, M.J.H. Ku, G. Roati, M.W. Zwierlein, Nature 472, 201 (2011)

Spin diffusion

 scattering conserves total + momentum: mass current preserved but changes relative + momentum: spin current decays



Is there a quantum limit for diffusion?



cf. spin Coulomb drag in GaAs quantum wells: $D_s\simeq 500\,\hbar/m$

Weber et al. 2005

- kinetic theory: diffusion coefficient $D_s \approx v \ell_{\rm mfp}, \ v \simeq \hbar k_F/m, \ \ell_{\rm mfp} \gtrsim 1/k_F$ $D_s \simeq \frac{\hbar}{-}$ quantum limit of diffusion
- spin conductivity from current correlations:

$$\sigma_s(\omega) = \frac{1}{\omega} \operatorname{Re} \int_0^\infty dt \, e^{i\omega t} \int d^3x \, \left\langle \left[j_s^z(\boldsymbol{x}, t), j_s^z(0, 0) \right] \right\rangle$$

with spin current operator $j_s(\boldsymbol{x},t) = j_{\uparrow}(\boldsymbol{x},t) - j_{\downarrow}(\boldsymbol{x},t)$

Dynamical spin conductivity



• exact high-frequency tail Hofmann PRA 2011; Enss & Haussmann PRL 2012

$$\sigma_s(\omega \to \infty) = \frac{C}{3\pi (m\omega)^{3/2}}$$

Spin diffusivity

• obtain diffusivity from Einstein relation, $D_s = \frac{\sigma_s}{\chi_s}$



- Quantum Monte Carlo simulation for finite lattice: $D_s\gtrsim 0.8\frac{\hbar}{m}$ Wlazlowski et al. PRL 2013

Thermal expansion and transport Frank/Zw. '14

is there a gravity dual of the unitary Fermi gas ?

expansion $\alpha_p = \kappa_T \left(\frac{\partial p}{\partial T}\right)_V \rightarrow \gamma \kappa_T c_v$ Grüneisen parameter γ scale invariance implies $p = 2\epsilon/3 \rightarrow \gamma(T) = 2/3$ is universal therm. conductivity $j_{Q} = -\kappa \nabla_{r} T$ at $j_{p} \equiv 0$ Boltzmann equ. in a 1/N-expansion $\kappa = 1.89 \dots N^2 T / \lambda_T \sim T^{3/2}$ **Prandtl-number** $\Pr = \eta c_P / \kappa = 1$ if a gravity dual exists (Son 08') leading order in 1/N gives Pr = 0.630136... + O(1/N)

The unitary gas is a benchmark for many-body physics. It

• realizes a high-temperature fermionic superfluid below

 $T_c/T_F \simeq 0.16$ and a scale-invariant many-body problem with universal ratios $p/p_F = \xi_s \simeq 0.37$ or $S/N|_c \simeq 0.7 k_B$

- exhibits universal fermionic spectral functions $A(k/k_F, \varepsilon/\varepsilon_F)$ in k-resolved RF and no pronounced pseudogap above T_c
- is the most perfect non-relativistic fluid with η/s close to the KSS bound and quantum-limited spin-diffusion $D_s\simeq 1.3\,\hbar/m$

open problems

- unconventional pairing in imbalanced gases, FFLO,
- Fermi and Bose polarons, quantum impurity problems



• transport in the quantum critical regime, solitons, ...

